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Soft theorems: Symmetry & Geometry

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Based on work with:

Berean-Dutcher, Cheung, Derda, Helset

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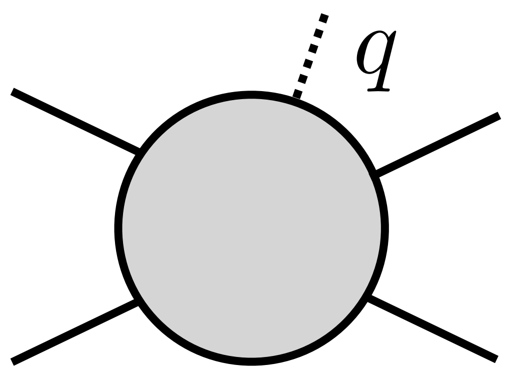
Soft theorems

- The behavior of scattering amplitudes when the momentum of a particle is small is often universal



$$\lim_{q \to 0} A_{n+1} = \mathcal{S} A_n$$

- Earliest example: Soft photon theorem [Low; Burnett, Kroll; Weinberg]

"leading" "subleading"

$$\lim_{q_\gamma \to 0} A_{n+1} = \sum_a \frac{q_a}{q \cdot p_a} \left[\epsilon \cdot p_a + \epsilon \cdot J_a \cdot q \right] A_n$$

similar for soft gluons, gravitons, pions

Perspectives on soft theorems

Symmetry

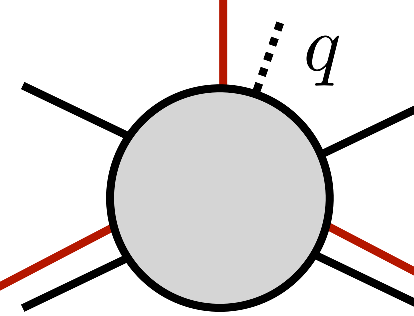
"Soft theorems are a consequence of symmetry"

e.g. asymptotic (photons, gravitons), spontaneously broken (pions)

Effective field theory

"Soft theorems are EFT of long-wavelenght modes"

e.g. factorization, ultrasoft decoupling in SCET



<u>Geometry</u>

This talk

Soft Nambu-Goldstone Bosons

- Spontaneous symmetry breaking implies existence of NGB

$$\pi^a \to \pi^a + v^a + \mathcal{O}(\pi^2) \qquad\qquad j^a_\mu = \partial_\mu \pi^a + \mathcal{O}(\pi^2)$$

- Also current algebra constrains their dynamics via soft theorems

$$\partial \cdot j^a(x) j^b_\mu(y) \sim f^{abc}\delta(x-y) j^c_\mu(y)$$

$$\Downarrow$$

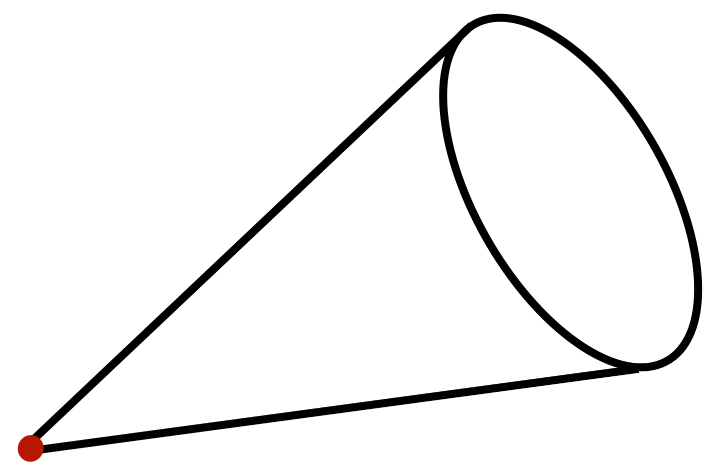
$$\lim_{q_\pi \to 0} A_{n+1} = 0$$ "Adler zero"

$$\lim_{q_{\pi^a}, q_{\pi^b} \to 0} A^{i_1 \cdots i_n i_a i_b}_{n+2} = \frac{1}{2} \sum_{c \neq a,b} \frac{s_{ac} - s_{bc}}{s_{ac} + s_{bc}} [\mathcal{X}^a, \mathcal{X}^b]^{i_c}{}_{j_c} A^{i_1 \cdots j_c \cdots i_n}_n$$

Soft moduli?

Many of our favorite theories have moduli spaces of vacua, parameterized by v.e.v of massless scalars

e.g., N=4 Coulomb branch



$\langle \Phi^I \rangle$

Protected by symmetry but <u>not always spontaneously broken</u> (e.g. SUSY moduli spaces)

Q: What is the general meaning of soft limits of scalar moduli?

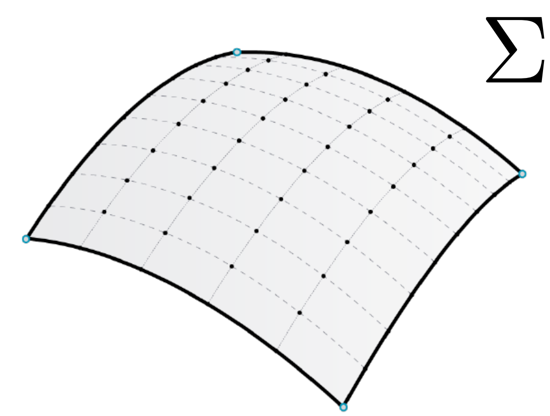
A: They encode the geometry of the moduli space!

Outline

- Amplitudes & geometry of moduli
- Geometric soft theorems
- Beyond scalars

Amplitudes & geometry of moduli

Review: Geometry of fields

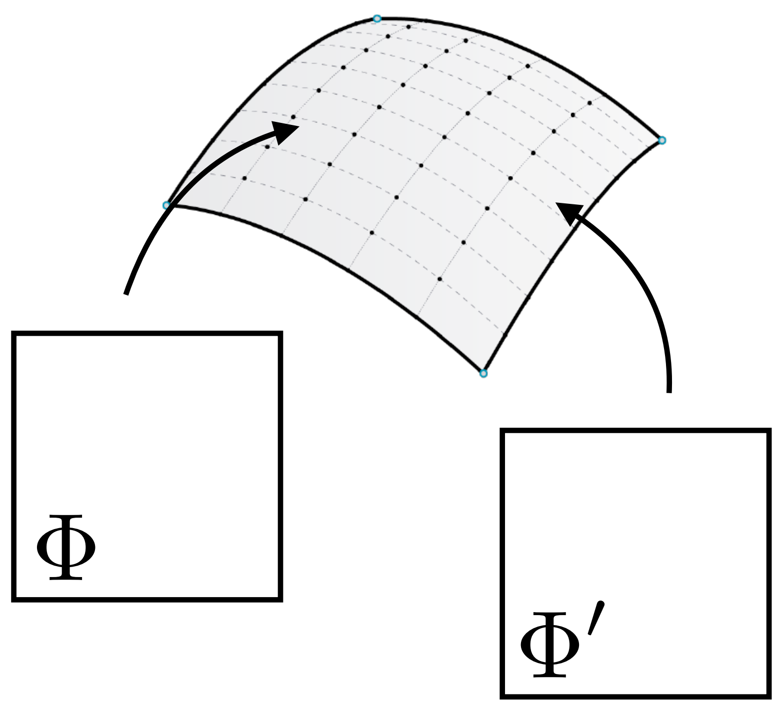


- Scalar fields take values in a target space manifold
- Lagrangian can be organized by derivative order

$$\frac{1}{2}g_{IJ}(\Phi)\partial_\mu\Phi^I\partial^\mu\Phi^J - V(\Phi) + \frac{1}{4}\lambda_{IJKL}(\Phi)\partial_\mu\Phi^I\partial^\mu\Phi^J\partial_\nu\Phi^K\partial^\nu\Phi^L + \cdots ,$$

- Field redefinitions = changes of coordinates $\Phi^I = \Phi^I(\Phi')$

$$\partial_\mu\Phi^I \quad\rightarrow\quad \frac{\partial\Phi'^I}{\partial\Phi^J}\partial_\mu\Phi^J$$



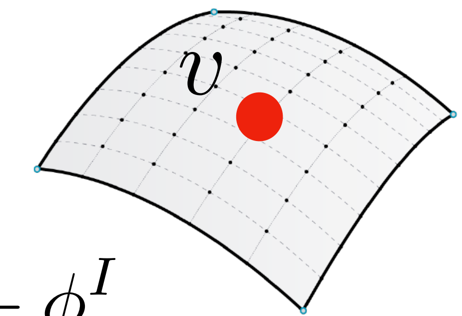
- Couplings are tensors in the target space

  e.g. two-derivative = metric

$$g_{IJ}(\Phi) \quad\rightarrow\quad \frac{\partial\Phi^K}{\partial\Phi'^I}\frac{\partial\Phi^L}{\partial\Phi'^J}g_{KL}(\Phi')$$

Familiar from world-sheet, but more general for EFT of moduli in D>2.

Geometry of amplitudes



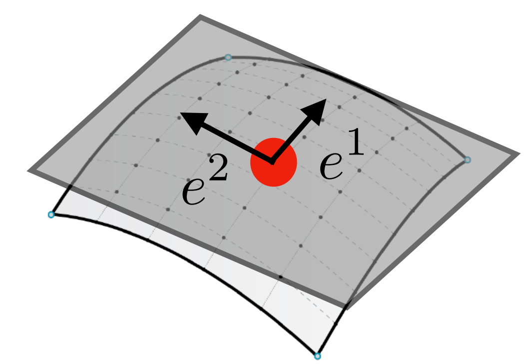
- Amplitudes defined by expanding around VEV $\Phi^I = v^I + \phi^I$

- Do not depend on field basis $\phi \to \phi + \epsilon f(\phi)$

$$S(\phi) \to S(\phi) + \frac{\delta S}{\delta \phi}\epsilon f(\phi) + \cdots$$

equations of motion

$$\propto p^2 - m^2 + \cdots$$



but on a choice of frame $\langle p^i | \phi^J(x) | 0 \rangle = e^{iJ}(v) e^{ip\cdot x}$

wavefunction ren.

- Must be a function of geometric invariants on $T\Sigma$!

e.g. curvature of metric connection on $T\Sigma$ $R^{ijkl}(v)$

[Volkov; Dixon, Kaplunovsky, Louis] $\nabla^m R^{ijkl}(v)$

Examples at tree level

- Two-derivative $\frac{1}{2} g_{IJ}(\Phi)\partial_\mu \Phi^I \partial^\mu \Phi^J$

$$
\begin{aligned}
A_4^{i_1 i_2 i_3 i_4} &= R^{i_1 i_3 i_2 i_4} s_{34} + R^{i_1 i_2 i_3 i_4} s_{24}, \\
A_5^{i_1 i_2 i_3 i_4 i_5} &= \nabla^{i_3} R^{i_1 i_4 i_2 i_5} s_{45} + \nabla^{i_4} R^{i_1 i_3 i_2 i_5} s_{35} + \nabla^{i_4} R^{i_1 i_2 i_3 i_5} s_{25} \\
&\quad + \nabla^{i_5} R^{i_1 i_3 i_2 i_4} s_{34} + \nabla^{i_5} R^{i_1 i_2 i_3 i_4} (s_{24} + s_{45}), \\
A_6^{i_1 i_2 i_3 i_4 i_5 i_6} &= -\tfrac{1}{72} \left(R^{i_1 i_3 i_2 j} s_{12} + R^{i_1 i_2 i_3 j} s_{13} \right) \frac{1}{s_{123}} \left(R_j{}^{i_6 i_5 i_4} s_{46} + R_j{}^{i_5 i_6 i_4} s_{45} \right) \\
&\quad + \tfrac{1}{108} \left(R^{i_1 i_3 i_2 j} (s_{12} - \tfrac{1}{6} s_{123}) + R^{i_1 i_2 i_3 j} (s_{13} - \tfrac{1}{6} s_{123}) \right) \left(R_j{}^{i_6 i_5 i_4} + R_j{}^{i_5 i_6 i_4} \right) \\
&\quad + \tfrac{1}{90} R^{i_1 i_6 i_5 j} R_j{}^{i_2 i_3 i_4} s_{13} + \tfrac{1}{80} \nabla^{i_6} \nabla^{i_5} R^{i_1 i_2 i_3 i_4} s_{13} + \text{perm}\,.
\end{aligned}
$$

- Four-derivative $\lambda_{IJKL}(\Phi)\partial_\mu \Phi^I \partial^\mu \Phi^J \partial_\nu \Phi^K \partial^\nu \Phi^L$

$$
\begin{aligned}
A_{4,\lambda}^{i_1 i_2 i_3 i_4} &= \tfrac{1}{2} s_{12} s_{34} \lambda^{i_1 i_2 i_3 i_4} + \tfrac{1}{2} s_{13} s_{24} \lambda^{i_1 i_3 i_2 i_4} + \tfrac{1}{2} s_{23} s_{14} \lambda^{i_2 i_3 i_1 i_4}\,, \\
A_{5,\lambda}^{i_1 i_2 i_3 i_4 i_5} &= \tfrac{1}{2} s_{12} s_{34} \nabla^{i_5} \lambda^{i_1 i_2 i_3 i_4} + \tfrac{1}{2} s_{13} s_{24} \nabla^{i_5} \lambda^{i_1 i_3 i_2 i_4} + \tfrac{1}{2} s_{23} s_{14} \nabla^{i_5} \lambda^{i_2 i_3 i_1 i_4} \\
&\quad + \tfrac{1}{2} s_{23} s_{45} \nabla^{i_1} \lambda^{i_2 i_3 i_4 i_5} + \tfrac{1}{2} s_{24} s_{35} \nabla^{i_1} \lambda^{i_2 i_4 i_3 i_5} + \tfrac{1}{2} s_{34} s_{25} \nabla^{i_1} \lambda^{i_3 i_4 i_2 i_5} \\
&\quad + \tfrac{1}{2} s_{13} s_{45} \nabla^{i_2} \lambda^{i_1 i_3 i_4 i_5} + \tfrac{1}{2} s_{14} s_{35} \nabla^{i_2} \lambda^{i_1 i_4 i_3 i_5} + \tfrac{1}{2} s_{34} s_{15} \nabla^{i_2} \lambda^{i_3 i_4 i_1 i_5} \\
&\quad + \tfrac{1}{2} s_{12} s_{45} \nabla^{i_3} \lambda^{i_1 i_2 i_4 i_5} + \tfrac{1}{2} s_{14} s_{25} \nabla^{i_3} \lambda^{i_1 i_4 i_2 i_5} + \tfrac{1}{2} s_{24} s_{15} \nabla^{i_3} \lambda^{i_2 i_4 i_1 i_5} \\
&\quad + \tfrac{1}{2} s_{12} s_{35} \nabla^{i_4} \lambda^{i_1 i_2 i_3 i_5} + \tfrac{1}{2} s_{13} s_{25} \nabla^{i_4} \lambda^{i_1 i_3 i_2 i_5} + \tfrac{1}{2} s_{23} s_{15} \nabla^{i_4} \lambda^{i_2 i_3 i_1 i_5}
\end{aligned}
$$

Geometric soft theorems

[Cheung, Helset, **JPM**]

New soft scalar theorem

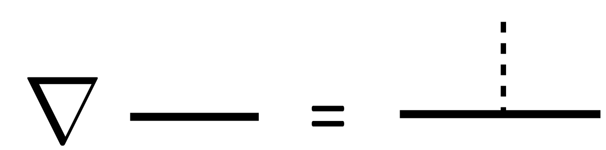
$$\lim_{q\to 0} A_{n+1}^{i_1\cdots i_n i} = \nabla^i A_n^{i_1\cdots i_n} + \sum_{a=1}^{n} \frac{\nabla^i V^{i_a}{}_{j_a}}{(p_a+q)^2 - m_{j_a}^2}\left(1 + q^\mu \frac{\partial}{\partial p_a^\mu}\right) A_n^{i_1\cdots j_a\cdots i_n}$$

subleading

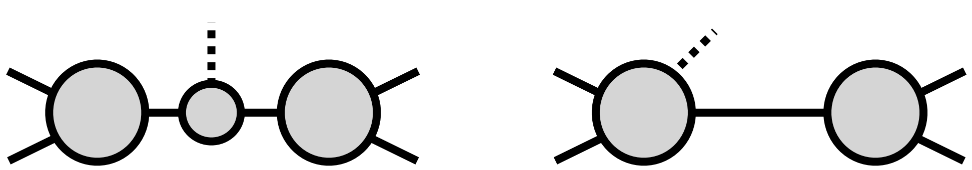
leading

Intuition:

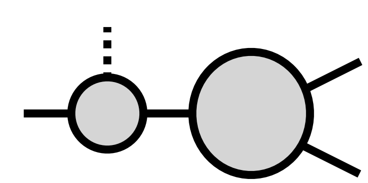
$$\lim_{q\to 0} A_{n+1} \sim \left(\nabla + \frac{\nabla m^2}{p^2 - m^2}\right) A_n$$



“Derivative w.r.t. VEV”



“on-shell connection”



Examples

- Two derivatives

$$A_4^{i_1 i_2 i_3 i_4} = R^{i_1 i_3 i_2 i_4} s_{34} + R^{i_1 i_2 i_3 i_4} s_{24},$$

$$\begin{aligned} A_5^{i_1 i_2 i_3 i_4 i_5} &= \nabla^{i_3} R^{i_1 i_4 i_2 i_5} s_{45} + \nabla^{i_4} R^{i_1 i_3 i_2 i_5} s_{35} + \nabla^{i_4} R^{i_1 i_2 i_3 i_5} s_{25} \\ &\quad + \nabla^{i_5} R^{i_1 i_3 i_2 i_4} s_{34} + \nabla^{i_5} R^{i_1 i_2 i_3 i_4} (s_{24} + s_{45}) \\ &= \nabla^{i_5} (R^{i_1 i_3 i_2 i_4} s_{34} + R^{i_1 i_2 i_3 i_4} s_{24}) \end{aligned}$$

- Four derivatives

$$\begin{aligned} A_{4,\lambda}^{i_1 i_2 i_3 i_4} &= \frac{1}{2} s_{12} s_{34} \lambda^{i_1 i_2 i_3 i_4} + \frac{1}{2} s_{13} s_{24} \lambda^{i_1 i_3 i_2 i_4} + \frac{1}{2} s_{23} s_{14} \lambda^{i_2 i_3 i_1 i_4} , \\ A_{5,\lambda}^{i_1 i_2 i_3 i_4 i_5} &= \frac{1}{2} s_{12} s_{34} \nabla^{i_5} \lambda^{i_1 i_2 i_3 i_4} + \frac{1}{2} s_{13} s_{24} \nabla^{i_5} \lambda^{i_1 i_3 i_2 i_4} + \frac{1}{2} s_{23} s_{14} \nabla^{i_5} \lambda^{i_2 i_3 i_1 i_4} \\ &\quad + \frac{1}{2} s_{23} s_{45} \nabla^{i_1} \lambda^{i_2 i_3 i_4 i_5} + \frac{1}{2} s_{24} s_{35} \nabla^{i_1} \lambda^{i_2 i_4 i_3 i_5} + \frac{1}{2} s_{34} s_{25} \nabla^{i_1} \lambda^{i_3 i_4 i_2 i_5} \\ &\quad + \frac{1}{2} s_{13} s_{45} \nabla^{i_2} \lambda^{i_1 i_3 i_4 i_5} + \frac{1}{2} s_{14} s_{35} \nabla^{i_2} \lambda^{i_1 i_4 i_3 i_5} + \frac{1}{2} s_{34} s_{15} \nabla^{i_2} \lambda^{i_3 i_4 i_1 i_5} \\ &\quad + \frac{1}{2} s_{12} s_{45} \nabla^{i_3} \lambda^{i_1 i_2 i_4 i_5} + \frac{1}{2} s_{14} s_{25} \nabla^{i_3} \lambda^{i_1 i_4 i_2 i_5} + \frac{1}{2} s_{24} s_{15} \nabla^{i_3} \lambda^{i_2 i_4 i_1 i_5} \\ &\quad + \frac{1}{2} s_{12} s_{35} \nabla^{i_4} \lambda^{i_1 i_2 i_3 i_5} + \frac{1}{2} s_{13} s_{25} \nabla^{i_4} \lambda^{i_1 i_3 i_2 i_5} + \frac{1}{2} s_{23} s_{15} \nabla^{i_4} \lambda^{i_2 i_3 i_1 i_5} \end{aligned}$$

Examples

- Two derivatives

$$
\begin{aligned}
A_4^{i_1 i_2 i_3 i_4} &= R^{i_1 i_3 i_2 i_4} s_{34} + R^{i_1 i_2 i_3 i_4} s_{24}, \\
A_5^{i_1 i_2 i_3 i_4 i_5} &= \nabla^{i_3} R^{i_1 i_4 i_2 i_5} \cancelto{0}{s_{45}} + \nabla^{i_4} R^{i_1 i_3 i_2 i_5} \cancelto{0}{s_{35}} + \nabla^{i_4} R^{i_1 i_2 i_3 i_5} \cancelto{0}{s_{25}} \\
&\quad + \nabla^{i_5} R^{i_1 i_3 i_2 i_4} s_{34} + \nabla^{i_5} R^{i_1 i_2 i_3 i_4} (s_{24} + \cancelto{0}{s_{45}}) \\
&= \nabla^{i_5} (R^{i_1 i_3 i_2 i_4} s_{34} + R^{i_1 i_2 i_3 i_4} s_{24})
\end{aligned}
$$

- Four derivatives

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\begin{aligned}
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&\quad + \cancelto{0}{\tfrac{1}{2} s_{23} s_{45} \nabla^{i_1} \lambda^{i_2 i_3 i_4 i_5} + \tfrac{1}{2} s_{24} s_{35} \nabla^{i_1} \lambda^{i_2 i_4 i_3 i_5} + \tfrac{1}{2} s_{34} s_{25} \nabla^{i_1} \lambda^{i_3 i_4 i_2 i_5}} \\
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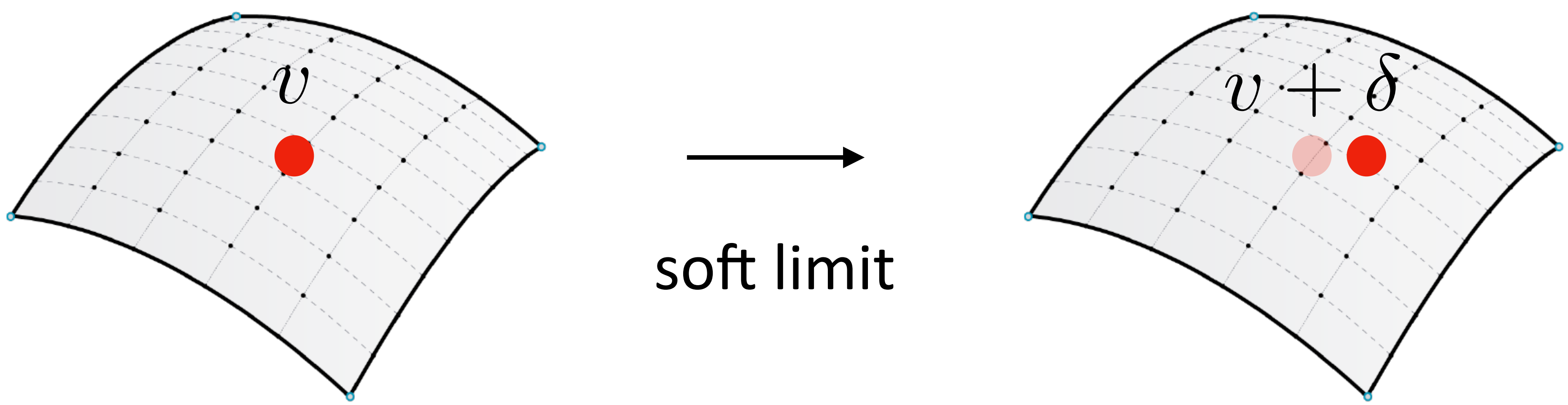
\end{aligned}
$$

Some comments

- Precise encoding of the intuition: “Soft scalar = shift of VEV”



- Proof is simple by treating VEV as a spurion
- Valid for **any** massless scalar, not only NGB
- Challenges common lore = soft scalar theorems *iff* SSB
- No symmetry required! soft theorems move us around space of vacua instead.

(Celestial interpretation in: [Kapec, Law, Narayanan])

Double soft measures curvature

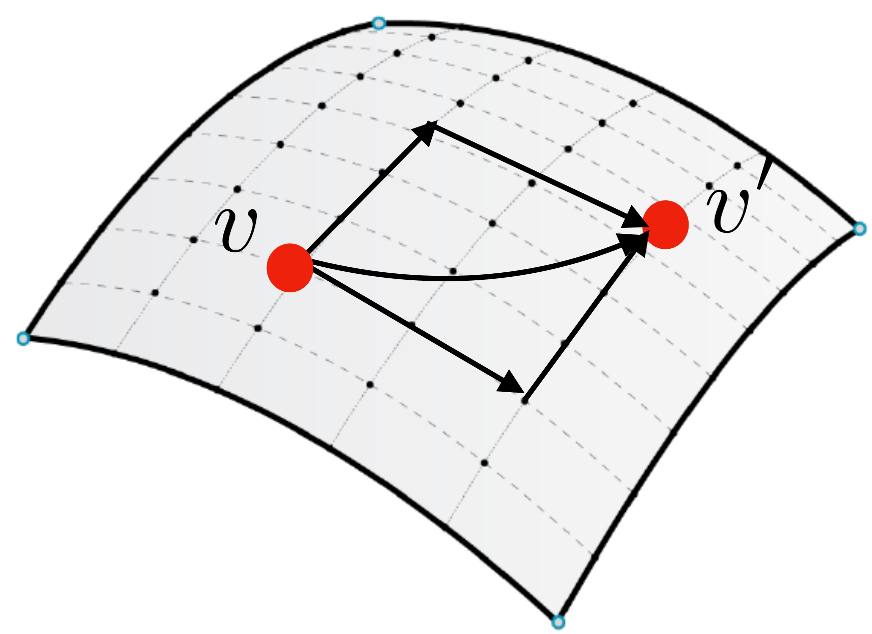
- Consecutive double soft

$$\left[\lim_{q_a\to 0}, \lim_{q_b\to 0}\right] A_{n+2}^{i_1\cdots i_n i_a i_b} = [\nabla^{i_a}, \nabla^{i_b}] A_n^{i_1\cdots i_n} = \sum_{c\neq a,b} R^{i_a i_b i_c}{}_{j_c} A_n^{i_1\cdots j_c \cdots i_n}$$

- Simultaneous double soft

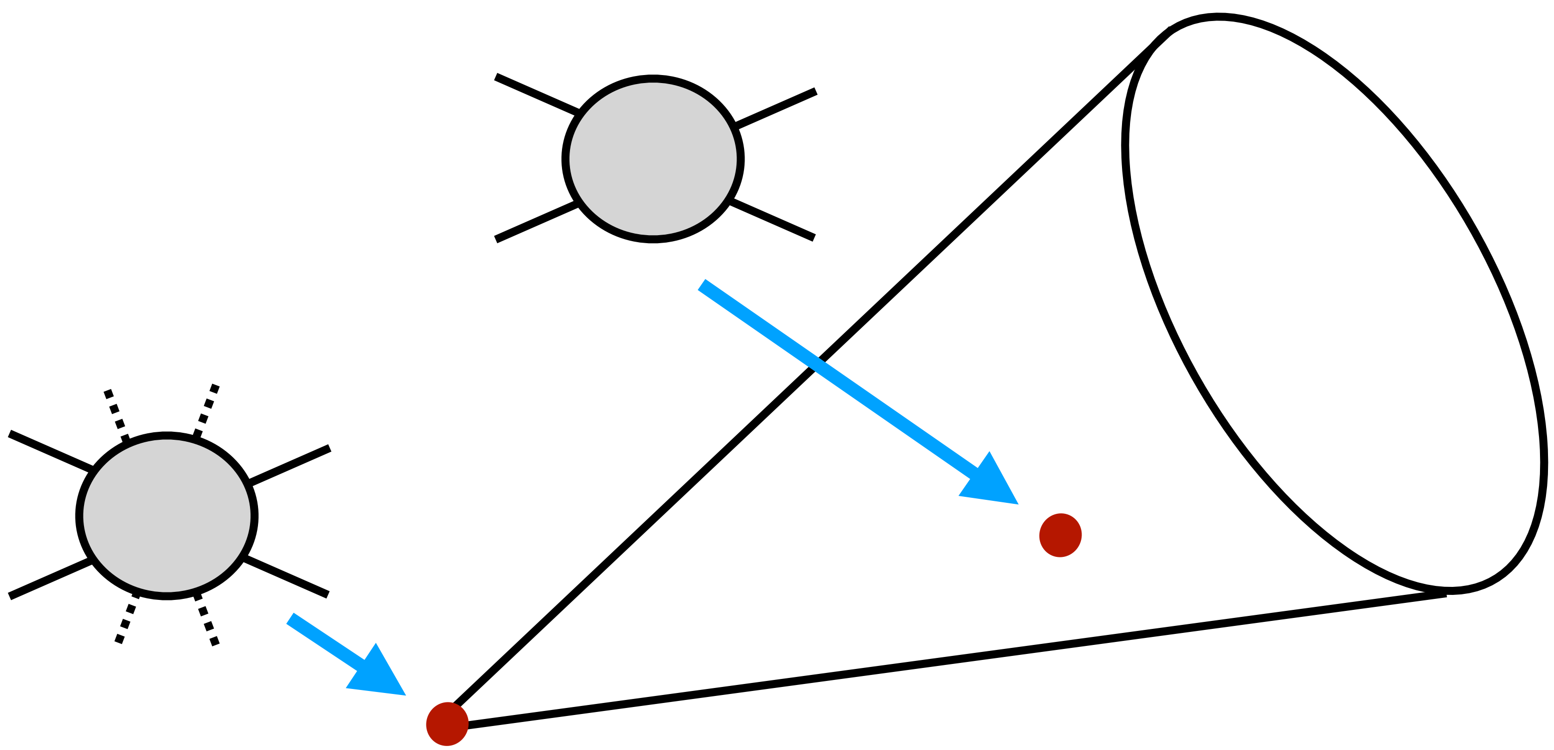
$$\lim_{q_a, q_b\to 0} A_{n+2}^{i_1\cdots i_n i_a i_b} = \frac{1}{2}\sum_{c\neq a,b} \frac{s_{ac}-s_{bc}}{s_{ac}+s_{bc}} R^{i_a i_b i_c}{}_{j_c} A_n^{i_1\cdots j_c\cdots i_n} + \nabla^{(i_a}\nabla^{i_b)} A_n^{i_1\cdots i_n}$$

- Difference: which path in field space



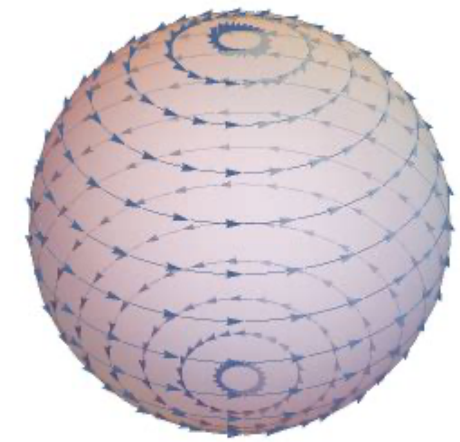
Exploring moduli space

- Geometric (multi)soft theorem lets us move around moduli space



- Requires knowledge of all amplitudes at given point in moduli space
- Converse enables soft recursion relations (ask me later)

Geometry of symmetry



- Symmetry = Killing vector

$$\Phi^I \quad \to \quad \Phi^I + \mathcal{K}^I(\Phi)$$

$$g_{IJ}(\Phi) \quad \to \quad g_{IJ}(\Phi) + \mathcal{L}_{\mathcal{K}} g_{IJ}(\Phi)$$

- Ward identity

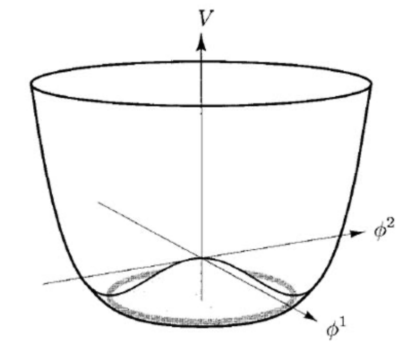
$$\mathcal{L}_{\mathcal{K}} A_n^{i_1 \cdots i_n} = \mathcal{K}_i \nabla^i A_n^{i_1 \cdots i_n} - \sum_{a=1}^{n} \nabla_{j_a} \mathcal{K}^{i_a} A_n^{i_1 \cdots j_a \cdots i_n} = 0$$

- Makes soft theorem multiplicative

$$\lim_{q \to 0} \mathcal{K}_i A_{n+1}^{i_1 \cdots i_n i} = \sum_{a=1}^{n} \nabla_{j_a} \mathcal{K}^{i_a} A_n^{i_1 \cdots j_a \cdots i_n}$$

Example: non-symmetric NBG

[Kampf, Novotny, Shifman, Trnka; Cheung, Helset, **JPM**]



- Spontaneously broken global symmetry

- Coset space G/H

$$[\mathcal{T}_a, \mathcal{T}_b] = f_{ab}{}^{c}\mathcal{T}_c\,,$$
$$[\mathcal{T}_a, \mathcal{X}_i] = f_{ai}{}^{j}\mathcal{X}_j\,,$$
$$[\mathcal{X}_i, \mathcal{X}_j] = f_{ij}{}^{a}\mathcal{T}_a + f_{ij}{}^{k}\mathcal{X}_k$$

- Soft theorem $\lim_{q\to 0} A_{n+1}^{i_1\ldots i_n i} = -\frac{1}{2}\sum_{a=1}^{n} f_{j_a}{}^{i_a i} A_n^{i_1\cdots j_a\cdots i_n}$

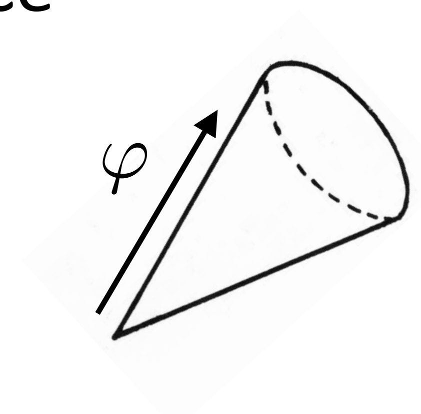
$\nabla\mathcal{X}$

- Symmetric coset ($\mathcal{X} \to -\mathcal{X}$) = Adler zero $\lim_{q\to 0} A_{n+1}^{i_1\ldots i_n i} = 0$

Example: soft dilaton

[Callan]

- Spontaneously broken spacetime symmetry - scale invariance
- Dilaton (NGB) = flat direction in moduli space $\nabla_\varphi = \partial_\varphi$



- New proof of soft dilaton theorem

$$\left(f_\varphi \partial_{\langle\varphi\rangle} + \sum_{a=1}^{n} p_a^\mu \frac{\partial}{\partial p_a^\mu} \right) A = (D - n\Delta) A$$

$$\lim_{q \to 0} A_{n+1} = \partial_{\langle\varphi\rangle} A_n = \frac{1}{f_\varphi} \left(D - n\Delta - \sum_{a=1}^{n} p_a^\mu \frac{\partial}{\partial p_a^\mu} \right) A_n$$

also works with masses!

Beyond scalars

Review: Coupling to matter

- Fermions specify vector bundle over moduli space $\mathcal{V}_f \to \Sigma$

$$h_{PQ}(\Phi)\bar{\psi}^P \gamma^\mu \overleftrightarrow{\partial}_\mu \psi^Q + \Omega_{PQI}(\Phi)\bar{\psi}^P \gamma^\mu \psi^Q \partial_\mu \Phi^I$$

fiber metric on $\mathcal{V}_f$ connection on $\mathcal{V}_f$

- Similarly, vectors live in $\mathcal{V}_v \to \Sigma$ with a metric $\frac{1}{2} h'_{AB}(\Phi) F^A \wedge \star F^B$

- Supersymmetry, among other things, can imply $\mathcal{V}_f \sim \mathcal{V}_v \sim T\Sigma$

Do soft moduli realize this geometry?

Soft scalars with matter

- Soft theorem works!

$$\lim_{q\to 0} A_{n+1} = \bar{\nabla} A_n = (\nabla + \Omega + \nabla h)\, A_n$$

connection on $T\Sigma \oplus \mathcal{V}_f \oplus \mathcal{V}_v$

- Charged scalars more involved: gives a version of Goldstone boson equivalence theorem.

- E.g. dipole coupling in D=4 $\quad D_{PQA}(\Phi)(\bar{\psi}^P \sigma^{\mu\nu} \psi^Q) W^A_{\mu\nu}$

$$\begin{aligned} A_3^{pqa} &= \langle 13\rangle\langle 23\rangle D^{pqa}, \\ A_4^{pqai_4} &= \langle 13\rangle\langle 23\rangle \bar{\nabla}^{i_i} D^{pqa}, \end{aligned}$$

Photon as NG boson

- Modern perspective: photon NG boson for $U(1)_e^{(1)} \times U(1)_m^{(1)}$

  e.g. electric 1-form symmetry

$$A_\mu \to A_\mu + \bar{A}_\mu \qquad\qquad \bar{A}_\mu \quad \text{Flat connection}$$

- Just like for pions, conserved current interpolates photon

$$J^{(1)}_{\mu\nu} = \partial_{[\mu} A_{\nu]} + \mathcal{O}(A^2)$$

- Symmetry is broken by charged matter

Soft photon

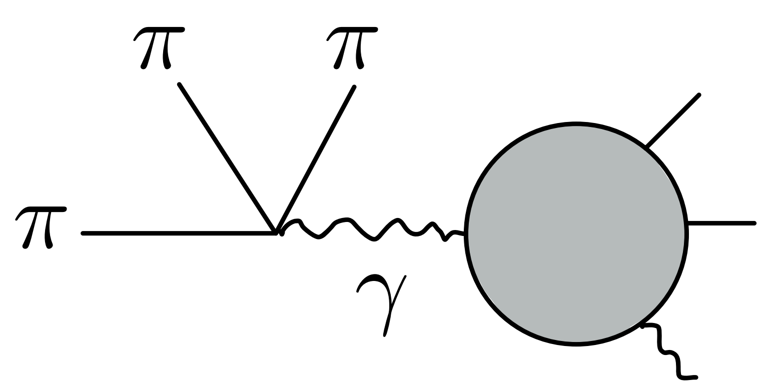
- In the absence of charged matter there is indeed an "Adler zero"

$$\lim_{q_\gamma \to 0} A_{n+1} = 0$$

- A bit boring, because higher-form symmetries abelian, and often only emergent

- Double soft gets interesting for two-group

$$\partial \cdot j^a(x) j^b_\mu(y) \sim f^{abc} \delta(x-y) j^c_\mu(y) + f^{ab\gamma} \partial^\mu \delta(x-y) J^{(1)}_{\mu\nu}(y)$$



Soft photon

- In the presence of charged matter

$$\lim_{q_\gamma \to 0} A_{n+1} = \sum_a \frac{q_a}{q \cdot p_a} \left[\epsilon \cdot p_a + \epsilon \cdot J_a \cdot q \right] A_n$$

- Both masslessness of photon and usual soft photon theorem stems from robustness of one-form symmetry

$$\delta \mathcal{L} = \bar{A}_\mu j^\mu \qquad\qquad \partial^\mu J^{(1)}_{\mu\nu} = j_\nu$$

- Intuition: Space of vacua still parameterized by $\bar{A}_\mu$

Geometric soft photon

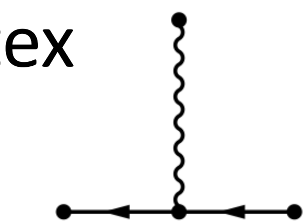
- Geometric proof with $\bar{A}_\mu$ as a spurion yields

$$\lim_{q\to 0} A_{n+1} \sim \left(\nabla + \bar{\Gamma}\right) A_n$$

$D = \partial + q\bar{A}$

$$\epsilon^\mu \frac{\delta}{\delta \bar{A}^\mu} = -\sum_a q_a \epsilon \cdot \frac{\partial}{\partial p_a}$$

Three-point vertex



$$q_a \frac{\epsilon\cdot p_a}{q\cdot p_a}\left(1 + p_a \cdot \frac{\partial}{\partial q}\right)$$

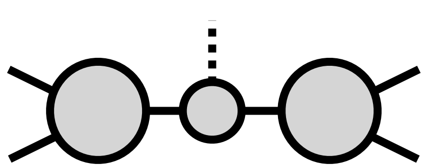
$$\sum_a \frac{q_a}{q\cdot p_a}\left[\epsilon \cdot p_a + \epsilon \cdot J_a \cdot q\right]$$ ✔

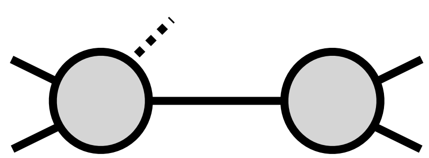
Note: space of vacua flat in this case

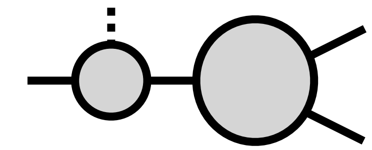
A general soft theorem?

Generalize the intuition:

$$\lim_{q\to 0} A_{n+1} \sim \left(\nabla + \bar{\Gamma}\right) A_n$$







"Soft particle = derivative w.r.t. flat background"

3pt-vertex "on-shell connection"

Does this work?

Conclusions

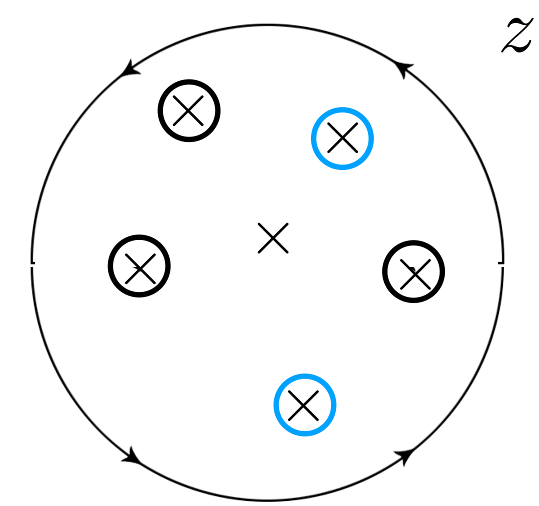
- Geometric perspective gives new and general soft theorems for scalar moduli beyond NG bosons
- Hints of general organizing principle for general soft theorems. Can we make this precise? Find new ones? Soft fermions beyond goldstino?
- Could enable new on-shell perspective of non-renormalization theorems
- Geometry perspective useful for SMEFT, soft recursion?
  [Helset, Trott, Alonso, Manohar, Jenkins, ….]
- Important question remains: Systematic way to move a finite distance in moduli space? Infinite? Massive amplitudes from massless amplitudes?

Thank you!

Geometric soft recursion

- Recursion with soft-subtractions

[Cheung, Kampf, Novotny, Shen, Trnka; Luo, Wen]



$$A_n(0) = \frac{1}{2\pi i}\oint \frac{dz}{z}\frac{A_n(z)}{F_{n,m}(z)} = -\sum_\alpha \mathrm{Res}_{z=z_\alpha^\pm}\left(\frac{A_n(z)}{zF_{n,m}(z)}\right)$$

$$F_{n,m}(z) = \prod_{a=1}^{n}(1-c_a z)^m$$

- Soft theorem enables on-shell recursion

[Cheung, Helset, **JPM**]

$$A_n(0) = \sum_\alpha \frac{A_L(z_\alpha^+)A_R(z_\alpha^+)}{(1-z_\alpha^+/z_\alpha^-)F_{n,1}(z_\alpha^+)} + (z_\alpha^+ \leftrightarrow z_\alpha^-) + \sum_a \frac{\nabla_{i_a}A_{n-1}(1/c_a)}{\Pi_{b\neq a}(1-c_b/c_a)}$$

- Seed of the recursion is four-point amplitude at every VEV