

Surface defects, phase transitions and supersymmetric black holes

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Motivation and overview

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- To quantitatively see this correspondence, we can consider the $\frac{1}{16}$ -BPS superconformal index, which can be computed faithfully in the weak-coupling regime via localization. The index counts the $\frac{1}{16}$ -BPS states, which matches the supersymmetric black hole entropy.

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- The superconformal index could be modified by inserting the extended observable. The extended observable should be at least $\frac{1}{16}$ -BPS and extend along temporal direction in order to probe the black hole geometry. The modified index reflects the entropy of (BH+string/brane) system.

Gukov-Witten surface defect in 4d SYM

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- This 4d-2d system has a holographic dual: Single surface defect insertion \longleftrightarrow Single probe D3 brane extending into the bulk and intersecting with the stack of N D3 branes (Constable, Erdmenger, Guralnik, Kirsch)

	0	1	2	3	4	5	6	7	8	9
N D3-branes	×	×	×	×						
Probe D3-brane	×	×			×	×				

Table: Bulk branes configuration

Superconformal index with surface defects

- The Hilbert space over the spatial S^3 is modified by including local operators on the defect: $\mathcal{H}(S^3) \rightarrow \mathcal{H}_{\mathcal{D}}(S^3)$. The defect index can be defined as a supersymmetric partition function tracing over the defect Hilbert space $\mathcal{H}_{\mathcal{D}}(S^3)$:

$$\mathcal{I}_{\mathcal{D}} = \text{Tr}_{\mathcal{H}_{\mathcal{D}}(S^3)} \left((-1)^F e^{-\beta\{\mathcal{Q}, \mathcal{Q}^\dagger\}} p^{J_1 + \frac{1}{2}R_3} q^{J_2 + \frac{1}{2}R_3} y_1^{\frac{1}{2}(R_1 - R_3)} y_2^{\frac{1}{2}(R_2 - R_3)} \right).$$

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- The index calculation can be reduced to a unitary matrix integral (Gadde, Gukov; Nakayama; Drukker, Gomis, Matsuura):

$$\mathcal{I}_{\mathcal{D}}(p, q, y_i) = \int_{SU(N)} [dU] \mathcal{I}_{4d}(p, q, y_i; U) \mathcal{I}_{2d}(p, q, y_i; U),$$

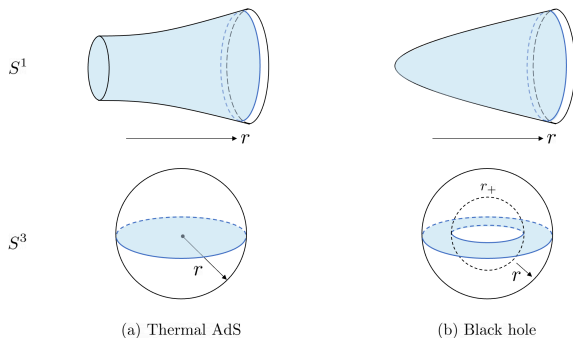
\mathcal{I}_{4d} and \mathcal{I}_{2d} are minimally coupled via $SU(N)$ gauge fields.

Large N saddle-point approximation (U_* is the saddle w.r.t. \mathcal{I}_{4d}):

$$\mathcal{I}_{\mathcal{D}}(p, q, y_i) = \sum_{\text{saddles } U_*} \mathcal{I}_{4d}(p, q, y_i; U_*) \mathcal{I}_{2d}(p, q, y_i; U_*) + \dots$$

Dual D3 brane action

- Probe D3 brane wraps the AdS_2 cigar geometry in BH phase.



- Evaluate the usual DBI action

$$I_{D3,E} = -T_{D3} \int \left(d^4x \sqrt{-\det(h_{D3})} - P[C_{(4)}] \right), \quad T_{D3} = \mathcal{O}(N)$$

$$\langle \mathcal{D} \rangle = \mathcal{I}_{2d}(p, q, y_i; U_*) \sim e^{-I_{D3,E}}.$$

Summary and Future direction

Conclusion

- Found Gukov-Witten surface defect as a supersymmetric order parameter of deconfinement phase transition: $\langle \mathcal{D} \rangle_{\text{thermal-AdS}} = e^{\mathcal{O}(1)}$,
$$\langle \mathcal{D} \rangle_{\text{BH}} = \exp \left\{ 2\pi i N \frac{(\sigma + \tau - 1)^2}{9\sigma} \right\}$$
- The holographic dual D3 brane, which interacts with the black hole, probes finer features of black hole physics

Future directions

- Other supersymmetric extended observables
- Half-BPS surface defects in other holographic theories
- Multiple surface defects insertion and backreaction to the geometry

Thank you!